

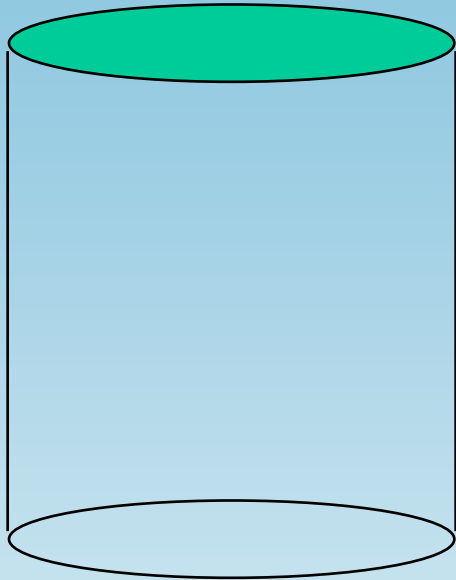
Gauss' Law for Cylinder Symmetry

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Presentations:

- Electromagnetism: History
- Electromagnetism: Electr. topics
- Electromagnetism: Magn. topics
- Electromagnetism: Waves topics
- Capacitor filling (complete)
- Capacitor filling (partial)
- Divergence Theorem
- E-field of a thin long charged wire
- E-field of a charged disk
- E-field of a dipole
- E-field of a line of dipoles
- E-field of a charged sphere
- E-field of a polarized object
- E-field: field energy
- Electromagnetism: integrations
- Electromagnetism: integration elements
- Gauss' Law for a cylindrical charge
- Gauss' Law for a charged plane
- Laplace's and Poisson's Law
- B-field of a thin long wire carrying a current
- B-field of a conducting charged sphere
- B-field of a homogeneously charged sphere

Gauss' Law for Cylinder Symmetry



Available:

Cylinder, radius R , infinitely long, carrying charge density λ [C/m]

Question:

Calculate E -field in arbitrary points inside and outside cylinder

Two cases:

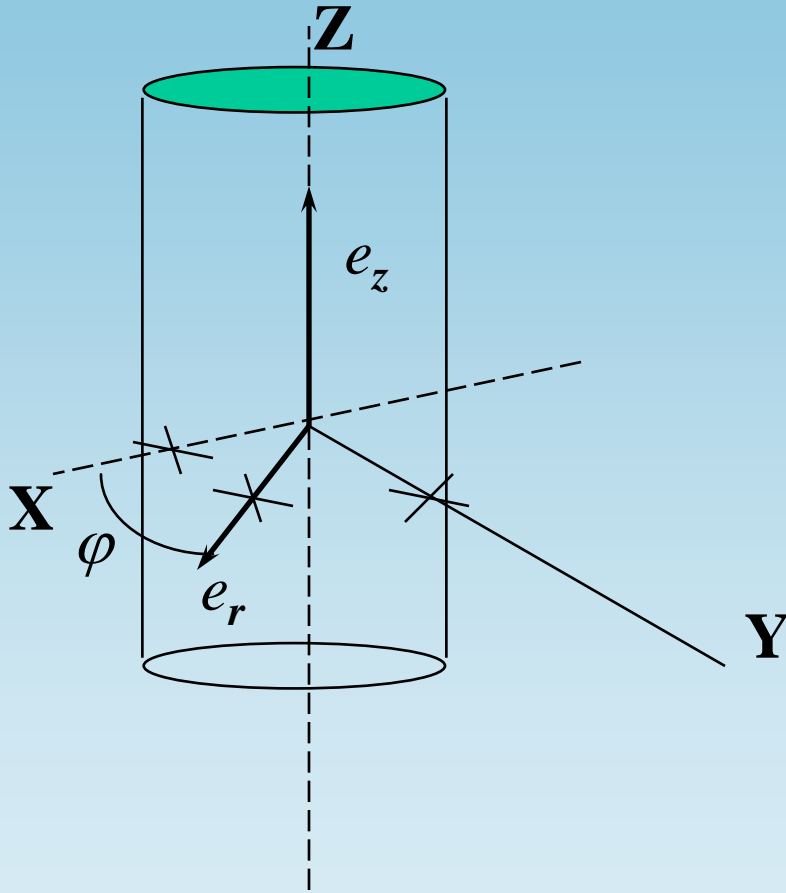
A: homogeneously charged

B: charged at surface walls only

Gauss' Law for Cylinder Symmetry

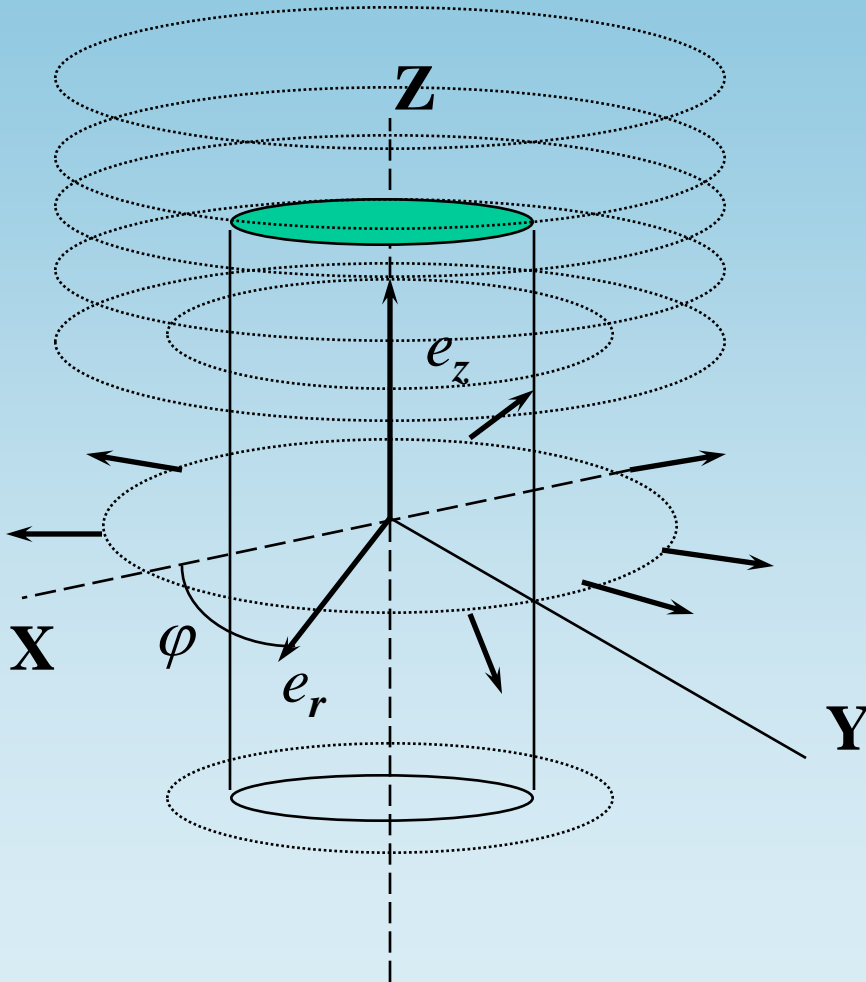
- Analysis and symmetry
- Approach to solution
- Calculations
- Conclusions

Analysis and Symmetry (1)



1. Cylinder: infinitely long, radius R
2. Charge distribution:
 λ [C/m] ; homogeneous.
3. Coordinate axes:
Z-axis = symm. axis
4. Cylinder symmetry:
all points at equal r are equivalent, even if at different z or φ

Analysis and Symmetry (2)



4. Cylinder symmetry:

all points at equal r are equivalent, even if at different z or φ .

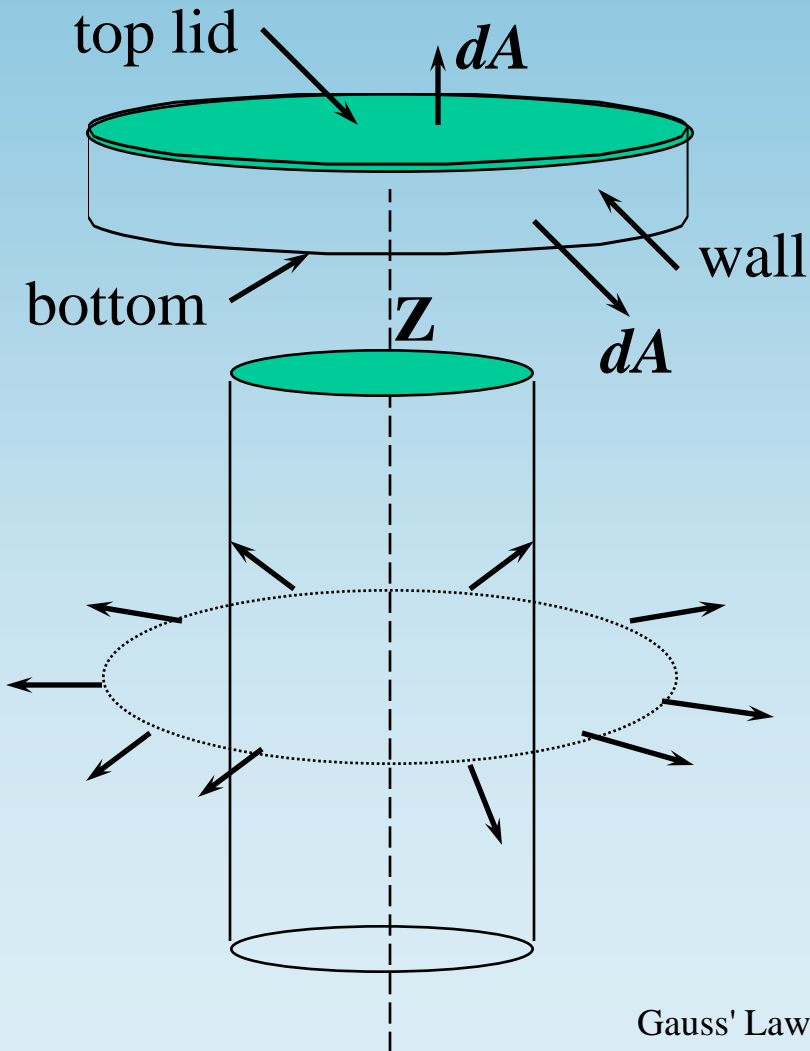
5. Consequences:

a point charge will not move tangentially.

E directed radially everywhere.

all planes $z = \text{const.}$ are equivalent.

Approach to solution



Gauss' Law:
$$\iint_A \mathbf{E} \cdot d\mathbf{A} = \frac{Q}{\epsilon_0}$$

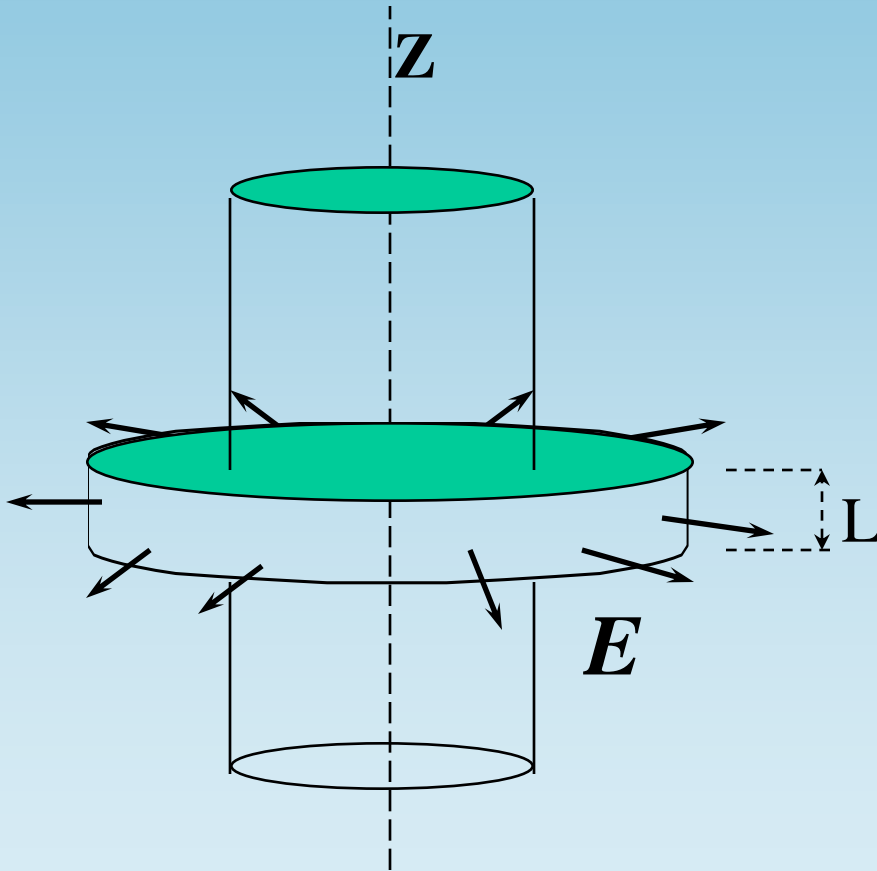
Choose Gauss-box A .

How to make optimum use of symmetry ??

- where $A \perp \mathbf{E}$
- where $A \parallel \mathbf{E}$
- where $E = 0$??

closed box needed !!
==> pill box

Calculations (1)



Gauss' Law:

$$\iint_A \mathbf{E} \cdot d\mathbf{A} = \frac{Q}{\epsilon_0}$$

pill box: radius $r > R$; height L

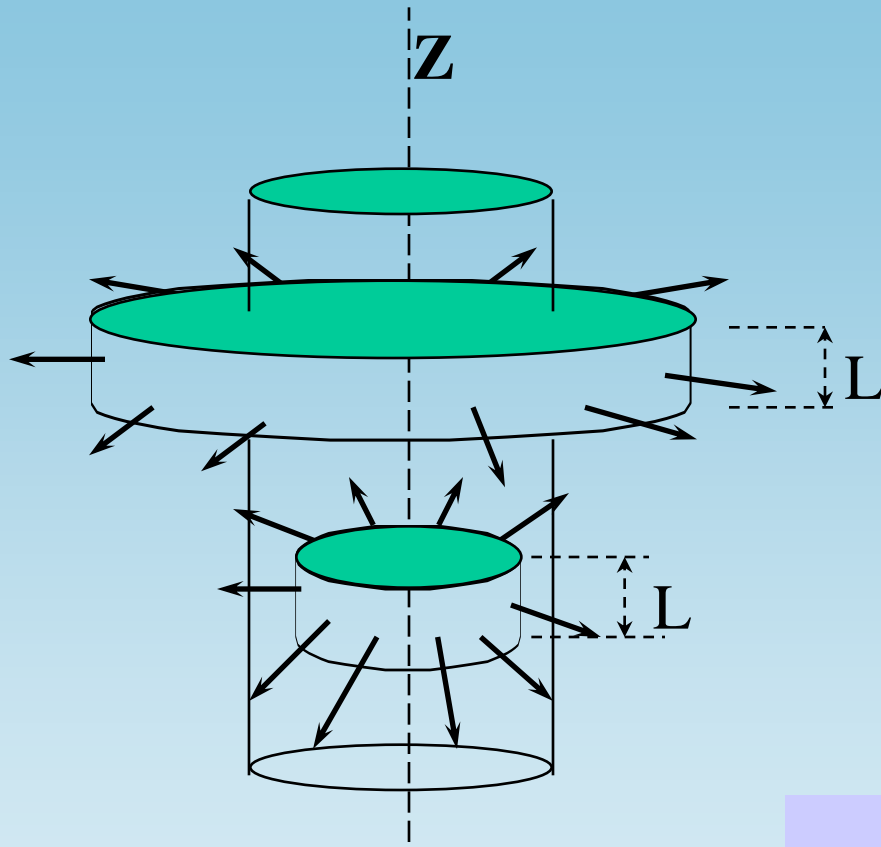
top and bottom lids do not
contribute ($\mathbf{E} \perp d\mathbf{A}$)

wall contributes: $E \cdot 2\pi r L$

charge enclosed: λL

result: $E(r) = \lambda / (2\pi\epsilon_0 r)$

Calculations (2)



A: homogeneously charged
B: charged at surface only

Gauss' Law:

$$\iint_A \mathbf{E} \cdot d\mathbf{A} = \frac{Q}{\epsilon_0}$$

result: $E(r) = \lambda / (2\pi\epsilon_0 r)$

but wait !! this holds for $r > R$!

for $r < R$: $\iint \Rightarrow E \cdot 2\pi r L$

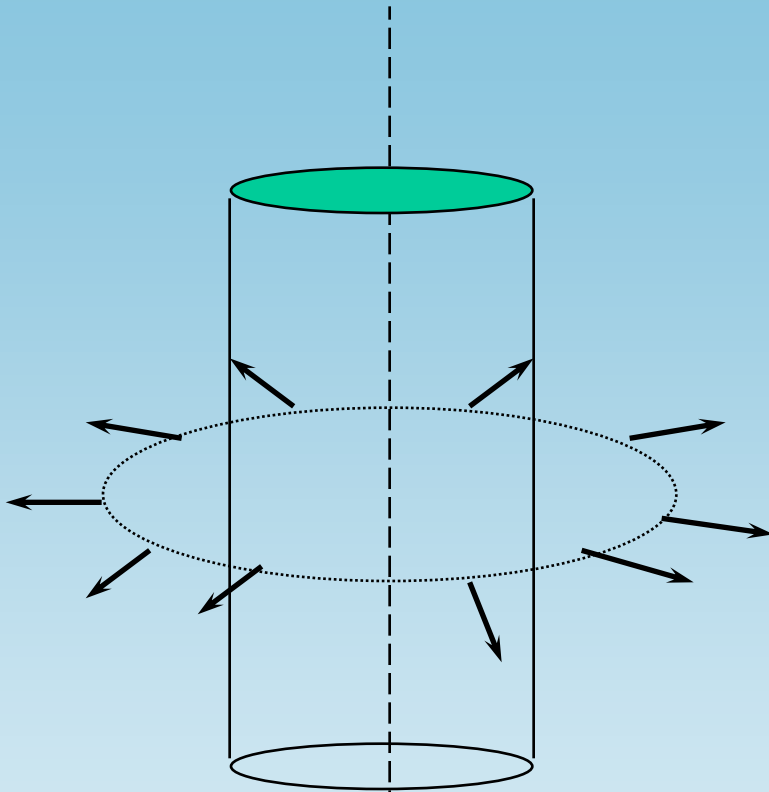
and $Q = Q(r) = \lambda L \frac{\pi r^2}{\pi R^2}$ (case A)

$Q = 0$ (case B)

$$\Rightarrow E(r) = \frac{\lambda L}{2\pi\epsilon_0 r L} \frac{r^2}{R^2} = \frac{\lambda r}{2\pi\epsilon_0 R^2} \quad (\text{case A})$$

$E(r) = 0$ (case B)

Conclusions (1)



A: homogeneously charged
B: charged at surface only

for infinite cylinder:

$$r > R: \mathbf{E}(r) = \frac{\lambda}{2\pi\epsilon_0 r} \mathbf{e}_r$$

$$r < R: \mathbf{E}(r) = \frac{\lambda r}{2\pi\epsilon_0 R^2} \mathbf{e}_r \quad (\text{case A})$$

$$\mathbf{E}(r) = 0 \quad (\text{case B})$$

field strength
dependent of distance
to cylinder =>
no homogeneous field

Conclusions (2)

for infinite cylinder: A: homogeneously charged;
B: surface charge only

$$r < R: \mathbf{E}(r) = \frac{\lambda r}{2\pi\epsilon_0 R^2} \mathbf{e}_r \text{ (A)}$$
$$E(r) = 0 \text{ (B)}$$

$$r > R: \mathbf{E}(r) = \frac{\lambda}{2\pi\epsilon_0 r} \mathbf{e}_r$$

